Multi-body systems

Wilhelm Kley Institut für Astronomie & Astrophysik & Kepler Center for Astro and Particle Physics Tübingen





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6. Multi-body systems:

Organisation

Lecture overview: 6.1 Resonances 6.2 Dynamics 6.3 Systems 6.4 Circumbinary planets

6.1 Resonances: Example from the Solar System

Orbital Periods: Neptune: 165yrs, Pluto: 248yrs Semi-major axis: Neptune: 29.7 AU, Pluto: 35.5 AU



Eccentricity $e_p = 0.25$ crossing orbits Stabilisation by: 3:2 Resonance Neptune: 3 Orbits Pluto: 2 Orbits Pluto shares orbits with plutinos \Rightarrow Pluto degraded from planet to dwarf planet Pluto-Flyby on 14.

July, 2015.

(NASA: New Horizon Mission)

6.1 Resonances: Example of orbits



(Gravity Simulator, Tony Dunn)

Planetary orbits: in the corotating-frame of Neptune Pluto Neptuns Uranus Saturn

Orbits avoid close encounters Pluto maximizes its separation from Neptune

The resonant condition stabilizes Pluto's orbit

6.1 Resonances: Definition - Mean motion resonance

A Mean Motion Resonance (MMR) is given by an integer ratio of their mean motions n_i

$$n_1/n_2 = p/q$$
 mit $q, p \in N$ (1)
orbital velocity of *i*-th planet (1: inner, 2: outer planet)

The resonant angles are then defined (with p > q)

$$\Phi_k = p\lambda_2 - q\lambda_1 - p\varpi_2 + q\varpi_1 + k(\varpi_2 - \varpi_1),$$
(2)

where λ_i and ϖ_i are the mean longitudes and longitudes of periapse. The integer *k* in Eq. (2) satisfies $q \le k \le p$ and has p - q + 1 possible values. Of the p - q + 1 resonant angles, at most two are linearly independent, and at least 1 angle will librate (resonant condition)

Example 2:1 System (p = 2 and q = 1)

n_i mean

$$\Phi_1 = 2\lambda_2 - \lambda_1 - \overline{\omega}_1, \quad \Phi_2 = 2\lambda_2 - \lambda_1 - \overline{\omega}_2 \qquad (\Delta \overline{\omega} = \Phi_1 - \Phi_2) \quad (3)$$

For details, see Book: Murray&Dermott Solar System Dynamics

- The resonant angles follow from an expansion of the perturbation potential between the two planets, that would move otherwise on unperturbed Kepler-orbits.
- Expand the perturbation in a Fourier-series
- The resonant angles are the stationary phases in the expansion
- If an angle Φ_i is in libration (i.e. covers a range smaller than [0, 2π]), then the system is said to be in a Resonance
- If both angles Φ_i (for a 2:1 resonance) librate (i.e. Δ_w as well) then the system is in apsidal corotation (ACR)

6.1 Resonances: Two sample systems

Animation of the observed motion of two systems

GJ 876: P₁=30d,P₂=60d



(Man Hoi Lee)

retrograde prcessionP $\Delta \varpi$ in libration $(\Delta \varpi)_{max} \approx 30^{\circ}$ Δ apsidal corotationDistance of periastrons: $\Delta \varpi = \varpi_2 - \varpi_1$

HD 128311: P1=458d,P2=928d



(Zsolt Sándor & Peter Klagyivik)

Planet 1: retrograde Precession $\Delta \varpi$ circulates

6.1 Resonances: Kepler-Planets



• but the systems lie slightly shifted from the resonance

6.1 Resonances: Formation of a resonant system

2 massive planets in disk (in HD 73526 \approx 2.5 $M_{\rm Jup}$ each)



6.1 Resonances: Eccentricity Pumping



6. Multi-body systems:

Lecture overview: 6.1 Resonances
6.2 Dynamics
6.3 Systems
6.4 Circumbinary planets

Organisation

6.2 Dynamics: Eccentricity vs. distance (RV-data)



6.2 Dynamics: Eccentricity and Inclination

What causes the high eccentricities and inclinations ?

In principle 2 models:

- differential migration of planets in the disk:
 - \Rightarrow compact (resonant) configuration
 - \Rightarrow disk dissipates and reduces the stabilisation
 - \Rightarrow instability of planetary system
 - \Rightarrow excitation of *e* and *i*
 - \Rightarrow gravitational scattering processes
- by an additional, massive body (star, planet): (*tidally driven migration*, by the Kozai-effect)
 - \Rightarrow third body excites eccentricity (inclination) to high values
 - \Rightarrow oscillation of eccentricity and inclination (of planet)
 - \Rightarrow for very high $e \Rightarrow$ tidal interaction with central star
 - \Rightarrow damping of semi-major axis

(not by disk-planet interaction)

6.2 Dynamics: Migration and scattering processes



N-body simulations with 3 or 4 planets

Start with:

- system near resonance

Consider:

- Star-Planet interaction

Black: Dots & circles Planets after 10⁹ yrs In brown - Kreise: Objekts after 10⁸ yrs In orange - squares: Obsevations

For more details: (Beauge & Nesvorny 2012) (Juvic & Tremaine 2008)

6.2 Dynamics: Kozai-Mechanism

A main-belt asteroid in the Solar System moves under the influence of the Sun and Jupiter. If the initial inclination of the body is higher than $i_0 > i_{crit} = 38^{\circ}$ the orbits will become highly eccentric (Kozai, 1962). Sometimes also Lidov-Kozai mechanism.



Consider a hierarchical restricted 3-body system

For asteroid:

$$m_0 = sun$$

$$m_1$$
 = asteroid

$$m_2 = \text{Jupiter}$$

for proto planet: $m_0 = \text{star}$ $m_1 = planet$ m_2 = secondary star

(Wen, 2003)

For the satellite, the orbit averaged equations of motion have a conserved quantity, $L_z = \cos(i)\sqrt{1-e^2}$, z is perpendicular to the orbit of the binary.

6.2 Dynamics: Kozai-Citations

Secular perturbations of asteroids with high inclination and eccentricity



6.2 Dynamics: The planetary system HD 80606



(Wikepedia)

HD 80606 a highly eccentric planetary system

$$P_p = 111 d$$

 $m_p = 3.94 M_{Jup}$
 $a_p = 0.449 AU$
 $e_p = 0.933$

6.2 Dynamics: Kozai mechanism for HD 80606

Consider the motion of the planet in the HD 80606 system under the action of an additional secondary star at large distance

Todays parameter: $m_p=3.94 \text{ M}_{Jup}$, $a_p = 0.449 \text{ AU}$, $e_p = 0.933$

start with: $a_p = 5 \text{ AU}$, $i = 85.6^\circ$, e = 0.1binary star parameter: $m_{bin} = 1.1 \text{ M}_{\odot}$, $a_{bin}=1000 \text{ AU}$, $e_{bin} = 0.5$

(e.g. Models by: Wu & Murray 2003, Fabrycky & Tremaine 2007)



6.2 Dynamics: Orbital elements of HD 80606



6. Multi-body systems: Organisation

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6.4 Circumbinary: Some observed Systems

Observed circumbinary planets (orbits normalized to the instability region)



is smaller than 2-5 times abin (Holman & Wiegert, 1999)

6.4 Circumbinary: Evolution of planets

What brings the planets so close to the instability limit ?

Assume planets have formed further out in pp-disk (flat systems) (not possible to form the in situ)

Study subsequent evolution of planets in disk (migration) Analyse two cases: Kepler-38 and Kepler-34 (See work by Pierens & Nelson (2013))



6.4 Circumbinary: Kepler 38: System Architecture



Binary Parameter: (Kepler 38A,B) $M_1 = 0.95 M_{\odot},$ $M_2 = 0.25 M_{\odot}$ $P_B = 18.6 \text{ days}$ $a_B = 0.15 \text{ AU}$ $e_B = 0.10$

Planet Parameter: (Kepler 38b) $m_p = 0.36 M_{Jup}$ $P_p = 105.6$ days $a_p = 0.46$ AU $e_p = 0.03$

6.4 Circumbinary: Kepler 38 - Circumbinary Disk Model

Two-dimensional disk around central binary, hydrodynamics, α -viscosity no self-gravity. Isothermal and radiative disks

Animation of disk around Binary Star (shown is surface density of inner disk)



Disk Parameter: Range: 0.25 - 2.0 AUMass: $3 \cdot 10^{-3} \text{ M}_{\odot}$ Sigma: 2000 g/cm² at 2 AU Equation of state: locally isothermal With central binary: truncates disk

(Kley & Haghighipour, 2014)

6.4 Circumbinary: Kepler 38 - Evolution of surface density



Kepler 38: Migration of planet in disk



Planets started at different distances, in original/relaxed disks All stop at same distance: at \approx 0.43 AU

Kepler 38: Surface density with planet

Planet is 'parked' at inner edge of disk



W. Kley Planet Formation, Multi-body systems, 45th Saas-Fee Lectures, 2015

Animation of Disk and Planet around Binary Star (shown is the final phase, in 5:1 resonance)



- 5:1 probable unstable upon disk dissipation
- some planet evolutions are unstable
- inside observed location and too high eccentricity
- vary disk properties to achieve stopping further out
- and use radiative disks

Results: Kepler 38:

- Circumbinary planets can be explained by a migration process
- Planets are 'parked' near gap edge, just outside of instability region
- Exact location depends on disk physical parameter

Open issues:

- Disk structure not clear even without planet
- Todo: Parameter variation, radiative disks, 3D simulations

Kepler 34: Eccentric binary, System Parameter

Comparison: Kepler 38 and Kepler 34 Binary Parameter: (Kepler 38) | (Kepler 34) $M_1[M_{\odot}]$ 0.95 1.05 $M_2[M_{\odot}]$ 0.25 1.02 P_B [days] 18.6 28 a_{B} [AU] 0.15 0.23 0.10 0.52 e_B Planet Parameter: (Kenler 38) (Kenler 34)

$m_p[M_{Jup}]$	0.36	0.21
P _p [days]	105.6	288
a _p [AU]	0.46	1.09
e _p	0.03	0.18

Kepler 34: Disk Structure

Surface density and eccentricity of disk with no planet



Differences to low eccentric binary:

- larger inner hole
- spike in density
- high disk eccentricity

Animation of Disk and Planet around Kepler 34 (shown is the final phase)



Summary: eccentric binary

- Large (eccentric) hole in center of disk
- Planet stops too far out
- Vary disk parameter ?
- Multiple planets ?

6. Multi-body systems: Organisation

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6.3 Systems: Systems of super-earths: Examples

Super-earth systems with $R_p < 1.5R_{earth}$ in comparison to the Solar System



Some characteristics of Hot-super Earth systems

 terrestrial-like exoplanets radii between Earth and Neptune often in multiple systems

extremely common 30-50% of Sun-line stars (FGK) host at least one planet with mass less than 10*M_{earth}* with period less than 50-100 d. Probability around M-stars at least as high

orbits

typically compact systems on non-resonant orbits orbital separations comparable to Solar System planets, in terms of Hill-radii

6.3 Systems: Systems of super-earths: Formation

Some suggested formation scenarios

- 1) in situ formation in massive disk
- 2) accretion during inward type I migration
- sheparding by interior mean-motion resonances of inwardly migrating massive planets
- sheparding by interior secular resonance inwardly migrating massive planets
- 5) circularisation of high-eccentric planets by tidal interactions with star
- 6) photo-evaporation of close-in giant planets

Numbers 3) to 6) can be ruled out theoretically and observationally (Raymond ea. 2014)

6.3 Systems: Systems of super-earths

Conflicts:

- 3)-4) require giant planet to drive planets inward and an additional disk to damp eccentricity in both scenarios a giant planet should be present just outside of outmost planet which is not obeserved
 - circularization in principle possible (see above) but require large initial eccentricities that would result in single-planet systems, which is in conflict with observations
 - 6) evaporation is possible for very close in planets, requires Gyr to operate. Strongly depends on distance from central star. Would only allow the innermost planet to be evaporated.

Hence 1) (in situ) and 2) (inward type I migration) are the leading contestors

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(Haghighiour 2013; Raymond ea. 2014)
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6.3 Systems: Systems of super-earths

Comments on scenarios

1) in situ

Requires migh mass ($\approx 20 - 40M_{\oplus}$) within fraction of an AU within a MMSN context: $\Sigma = \Sigma_0 (r/AU)^{-x}$, need steep powerlaw

x = 1.6 - 1.7 and 10 times higher normalisation Σ_0 .

i.e. high total disk mass

Observations indicate shallower profile: $x \approx 0.5 - 1.0$

But: Simulations show that orbital properties (eccentricities, inclinations, separations) match observations

2) migration

Simulations of migration super-earths show the formation of resonant chains

destabilisation \Rightarrow collisions and accretion

consecutive chain formation (see example below)

Possible distinction betwen models:

1) naked high-density rocks 2) lower density material containing ice but possible atmosphere could hide effect

6.3 Systems: Example formation scenario



Formation of hot super-Earths by type I migration Top: orbital distances Middle: mass growth Bottom: Comparision

size scaled to radius of the planet

For reference: $a_{Earth}/a_{Venus} = 1.6$

(Cossou ea. 2014)

Resonances:

Murray & Dermott (1999), Kley et al. (2004)

- Multiple: Fabrycky & Tremaine (2007)
- Super-Earths:

Haghighipour (2013), Raymond et al. (2014)